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Exercises for Computational Methods in Solid State Theory  $$\mathrm{SS}\ 2014$$ 

Exercise Set 4

(Due date: Monday, June 9, 2014)

## **Exercise 4** (Hubbard I approximation of the Anderson impurity model) (10 points)

We use equations of motion within a mean-field decoupling scheme to solve the Anderson impurity model

(1)

$$\mathcal{H} = \sum_{k\sigma} \varepsilon_{k} c_{k\sigma}^{\dagger} c_{k\sigma} + \sum_{\sigma} \varepsilon_{d} d_{\sigma}^{\dagger} d_{\sigma} + \frac{U}{2} \sum_{\sigma} \hat{n}_{\sigma} \hat{n}_{\bar{\sigma}} + \sum_{k\sigma} \left( V_{k\sigma}^{*} c_{k\sigma}^{\dagger} d_{\sigma} + V_{k\sigma} d_{\sigma}^{\dagger} c_{k\sigma} \right)$$

where  $\hat{n}_{\sigma} = d_{\sigma}^{\dagger} d_{\sigma}$ .

a) Use the general equation of motion for the Greens function

(2) 
$$\omega\langle\!\langle \mathbf{A};\mathbf{B}\rangle\!\rangle = \left\langle \left[\mathbf{A},\mathbf{B}\right]_{+}\right\rangle + \left\langle\!\langle \left[\mathbf{A},\hat{\mathbf{H}}\right]_{-};\mathbf{B}\right\rangle\!\rangle$$

to derive the following three equations:

(3) 
$$(\omega - \epsilon_{\mathbf{d}}) \langle\!\langle \mathbf{d}_{\sigma}; \mathbf{d}_{\sigma}^{\dagger} \rangle\!\rangle = 1 + U \langle\!\langle \mathbf{d}_{\sigma} \hat{\mathbf{n}}_{\bar{\sigma}}; \mathbf{d}_{\sigma}^{\dagger} \rangle\!\rangle + \sum_{\mathbf{k}} V_{\mathbf{k}\sigma} \langle\!\langle \mathbf{c}_{\mathbf{k}\sigma}; \mathbf{d}_{\sigma}^{\dagger} \rangle\!\rangle$$

(4) 
$$(\boldsymbol{\omega} - \boldsymbol{\varepsilon}_{\mathbf{k}}) \langle\!\langle \mathbf{c}_{\mathbf{k}\sigma}; \mathbf{d}_{\sigma}^{\dagger} \rangle\!\rangle = \mathbf{V}_{\mathbf{k}\sigma}^{*} \langle\!\langle \mathbf{d}_{\sigma}; \mathbf{d}_{\sigma}^{\dagger} \rangle\!\rangle$$

$$\begin{array}{lll} (5) \\ (\omega - \varepsilon_{d} - U) \langle\!\langle d_{\sigma} \hat{n}_{\bar{\sigma}}; d_{\sigma}^{\dagger} \rangle\!\rangle & = & \langle \hat{n}_{\bar{\sigma}} \rangle + \sum_{k} V_{k\bar{\sigma}}^{*} \langle\!\langle c_{k\bar{\sigma}}^{\dagger} d_{\sigma} d_{\bar{\sigma}}; d_{\sigma}^{\dagger} \rangle\!\rangle + \sum_{k} V_{k\sigma} \langle\!\langle c_{k\sigma} \hat{n}_{\bar{\sigma}}; d_{\sigma}^{\dagger} \rangle\!\rangle \\ & & - \sum_{k} V_{k\bar{\sigma}} \langle\!\langle c_{k\bar{\sigma}} d_{\bar{\sigma}}^{\dagger} d_{\sigma}; d_{\sigma}^{\dagger} \rangle\!\rangle \end{array}$$

b) Use the equations of motion derived in a) to show that the d electron Greens function can be written as

(6) 
$$\langle\!\langle \mathbf{d}_{\sigma}; \mathbf{d}_{\sigma}^{\dagger} \rangle\!\rangle = \frac{1 + \frac{\mathbf{U} \langle \hat{\mathbf{n}}_{\bar{\sigma}} \rangle}{\boldsymbol{\omega} - \boldsymbol{\varepsilon}_{\mathbf{d}} - \mathbf{U} - \Delta(\boldsymbol{\omega})}}{\boldsymbol{\omega} - \boldsymbol{\varepsilon}_{\mathbf{d}} - \Delta(\boldsymbol{\omega})}$$

where the hybridization function  $\Delta(\omega)$  is given as:

(7) 
$$\Delta(\omega) = \sum_{k} \frac{V^2}{\omega - \varepsilon_k}$$

For some of the four operator terms a mean-field decoupling is necessary to close the set of equations, e.g.

(8) 
$$\langle\!\langle \mathbf{c}_{\mathbf{k}\sigma}\mathbf{d}_{\bar{\sigma}}^{\dagger}\mathbf{d}_{\bar{\sigma}};\mathbf{d}_{\sigma}^{\dagger}\rangle\!\rangle \approx \left\langle\!\langle \mathbf{d}_{\bar{\sigma}}^{\dagger}\mathbf{d}_{\bar{\sigma}}\right\rangle \langle\!\langle \mathbf{c}_{\mathbf{k}\sigma};\mathbf{d}_{\sigma}^{\dagger}\rangle\!\rangle$$

## Exercise 5 (Anderson impurity model and Hubbard model on the Bethe lattice) (10 points)

We use the Hubbard I impurity solver derived in Exercise 4b) to solve the Anderson impurity model and dynamical mean-field theory (DMFT) for the Hubbard model on the Bethe lattice. All calculations should be done in the paramagnetic phase, where  $\langle n_{\sigma} \rangle = \langle n_{\bar{\sigma}} \rangle$ .

The Hubbard I solver can be used on the imaginary frequency axis by replacing  $\omega \to i\omega_n$ . For the calculation of occupation numbers on the Matsubara axis the high-frequency tails of the Greens function must be taken into account analytically.

The parameter  $\epsilon_d$  must be set to  $-\frac{U}{2}$  to make the impurity solver symmetric with respect to the Fermi energy. Mind that the lattice only enters through the density of states, which is semicircular for the Bethe lattice.

- a) Use the expression for the d electron Greens function to solve the Anderson impurity model on the Bethe lattice with a band width of W = 4 eV. Consider an energy independent hybridization of  $V = \sqrt{0.4}$  eV and a temperature T = 0.1 eV. Plot the density of states of the d electrons.
- b) Use the impurity solver of b) to solve the Hubbard model with dynamical mean field theory on the Bethe lattice (W, T as in a)). Find the critical interaction  $U_c$  for the metal to insulator transition at half filling. Plot the converged interacting density of states.

## Exercise 6 (DMFT solution of the one-band Hubbard model for a cuprate) (10 points)

Use dynamical mean field theory with the Hubbard I impurity solver to solve the one-band Hubbard model for the cuprate of Exercises 2 and 3. Here, U and T are U = 6 eV and T = 0.1 eV, respectively. Consider only the paramagnetic phase.

The parameter  $\epsilon_d$  must be adjusted in order to conserve the initial filling of the tight binding model.

Plot the converged interacting density of states.