Goethe-Universität Frankfurt Fachbereich Physik

Institut für Theoretische Physik Dr. Harald O. Jeschke Dr. Hunpyo Lee



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Exercises for Computational Methods in Solid State Theory SS 2012

Exercise Set 3

(Due date: Monday, May 14, 2012)

Exercise 3 (Matsubara Greens functions and Padé method) (10 points)

Consider the one-band tight binding model that you determined for $Sr_2CuO_2Cl_2$ in Exercise 2.

a) Calculate the noninteracting retarded Matsubara Greens functions

(1)
$$G(i\omega_n, \vec{k}) = \frac{1}{i\omega_n - \varepsilon(\vec{k})}$$

and determine the local Greens function

(2)
$$G_0(i\omega_n) = \sum_{\vec{k}} G(i\omega_n, \vec{k})$$

by explicitly performing the \vec{k} summation. $\varepsilon(\vec{k})$ is the tight binding dispersion. Consider a temperature $\beta = 40 \text{ eV}^{-1}$.

b) Use the Hilbert transformation which is given as

(3)
$$G_0(i\omega_n) = \int_{-\infty}^{\infty} d\varepsilon \frac{\rho(\varepsilon)}{i\omega_n - \varepsilon},$$

where the density of states $\rho(\epsilon)$ for the Cu $3d_{x^2-y^2}$ band crossing the Fermi level was determined by tight binding fit in Exercise 2. Calculate $G_0(i\omega_n)$ via Hilbert transformation and compare to $G_0(i\omega_n)$ obtained from a).

- c) Implement the method of Padé approximants and continue $G_0(i\omega_n)$ of Eq. (2) analytically to the real frequency axis. Plot the density of states $\rho(\omega) = -\frac{1}{\pi} \operatorname{Im} G_0(\omega)$ and compare to the tight binding result.
- d) Use your tight binding density of states $\rho(\omega)$ from Exercise 2 to determine $G_0(\omega)$ on the real frequency axis via Kramers-Kronig relations. Compare the result to $G_0(\omega)$ from c).